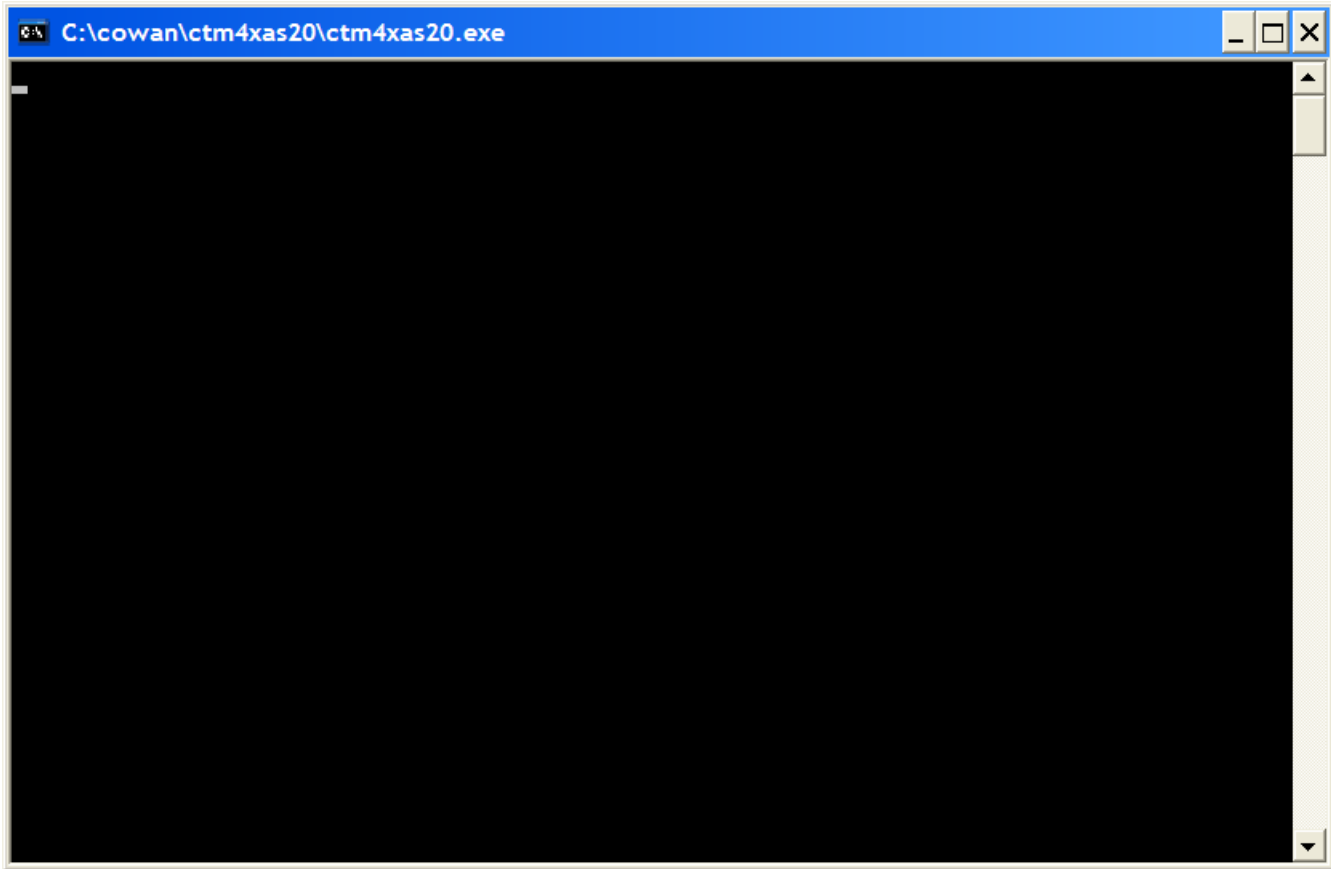


## Steps for successfully using the CTM4XAS20 Charge-Transfer Multiplet Program

Layout of inputs is different for the current CTM4XAS55 version, but this guide is still totally relevant and usable

- Go to <http://www.anorg.chem.uu.nl/people/staff/FrankdeGroot/SLs/ctm4xas2008.htm>. Follow the download and installation steps exactly. Finally, everything needed should be in the folder: c:\cowan\ctm4xas20\
  - Go to this folder and run ctm4xas20.exe
  - You will see a cmd window, as on slide 2, while the program gets ready to work
  - After a while, you will then see the real input page, as on slide 3, where various inputs and options are indicated
  - If you select charge transfer, which is optional for XAS since the excited electron is mostly excited into a 3d “screening” orbital, but required for XPS since the core hole will strongly polarize everything around it, other options appear, as shown on slide 4.
  - Background on the method and the parameters is shown in slides 5-7
  - If you now select XPS, you see something like slide 8
  - Now input some parameters appropriate to your case. For example, for Mn<sup>2+</sup> in MnO, with origins discussed in slides 9-12.



**Figures - CTM4XAS 2.0**

Calculate Plot Debug Desktop Window Help

**Yes or no for ligand-metal charge transfer**

**Don't use XAS, unless this is clicked**

Configuration

Ion   Auto Plot  XPS  Charge transfer

Initial state 2P06 3D08 Initial state (CT)

Final state 2P05 3D09 Final state (CT)

Slater integral reduction (%)

3d spin-orbit    Expert options

**Chooses whether a core 2p or 3p electron is excited**

Crystal field parameters (eV)

Symmetry     T(b1)

Initial state Final state

10 Dq     T(a1)

Dt     T(b2)

Ds     T(e)

M (meV)

**Simplest Octahedral symmetry, two others possible**

Large transfer parameters (eV)

Udd   T(a1)

U<sub>dd</sub> cannot be changed in XPS

U<sub>pd</sub>

**Shows XAS even when XPS plotted**

Plotting

File

Spectrum   Suppress sticks

Gaussian broadening

Lorentzian broadening

range (eV)

Force energy range  Split

**Click here to get rid of a file if you don't want to plot it**

**Don't force energy range, program selects it for plot**

**Lorentzian (lifetime) & Gaussian (instrumental) broadenings to simulate a true experimental spectrum. Set to zero to see details of states involved.**

Use the curled arrow to enlarge for full viewing of all inputs. There is some overlap of inputs on my screen at the bottom, but maybe not yours; no big problem anyway.

**CTM4XAS credits**

**CTM4XAS**

Charge Transfer Modeling Calculations for X-Ray Absorption Spectroscopy

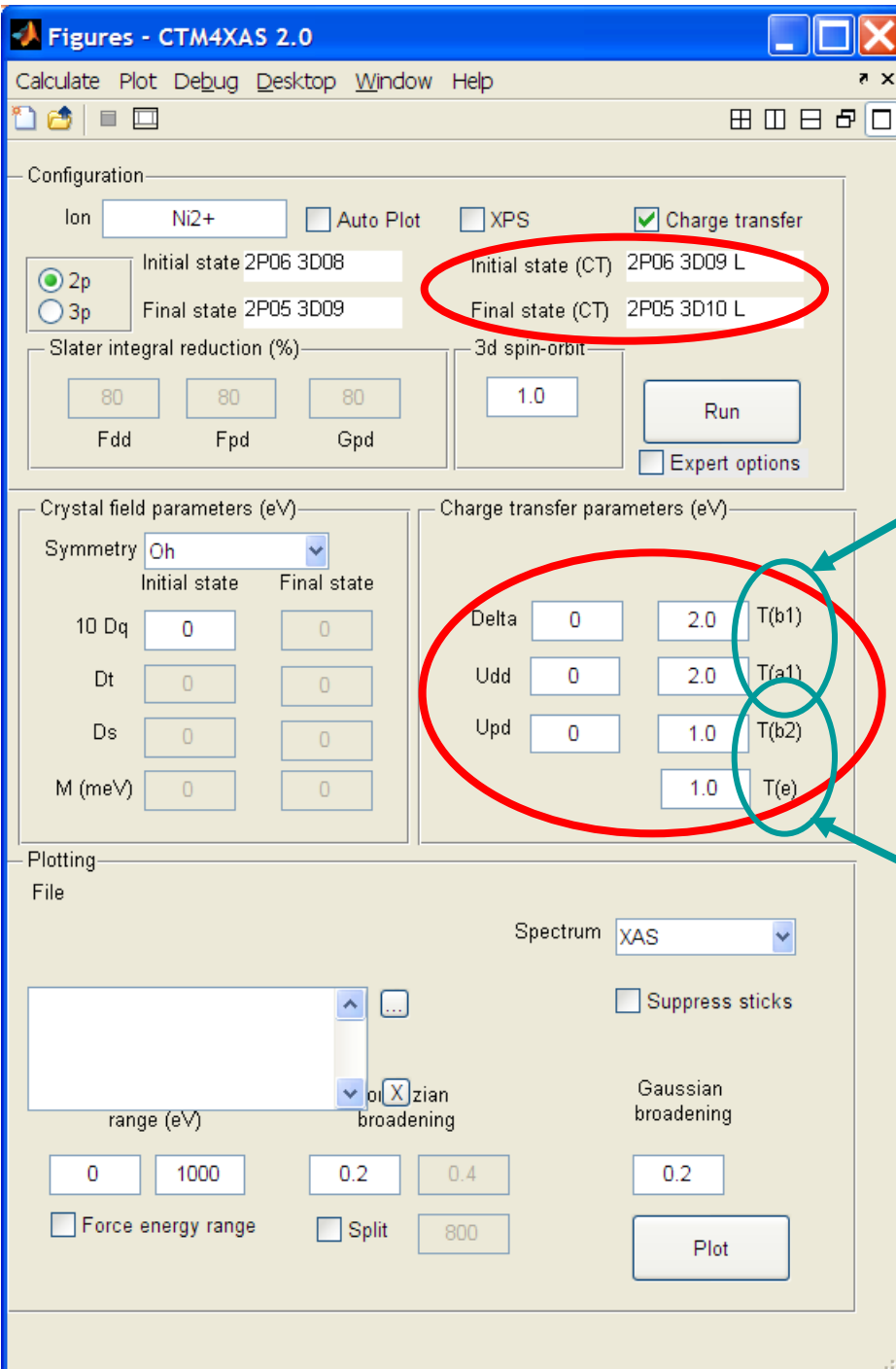
**NOW WITH XPS**

Version 2.0

© Eli Stavitski and Frank de Groot, 2008  
Inorganic Chemistry & Catalysis, Utrecht University

Just delete this advertisement!

Lorentzian (lifetime) & Gaussian (instrumental) broadenings to simulate a true experimental spectrum. Set to zero to see details of states involved.



## And with charge transfer selected

Initial and final states in XAS have some total number of electrons, and screened  $d^{n+1}$  configurations mixed in are now present

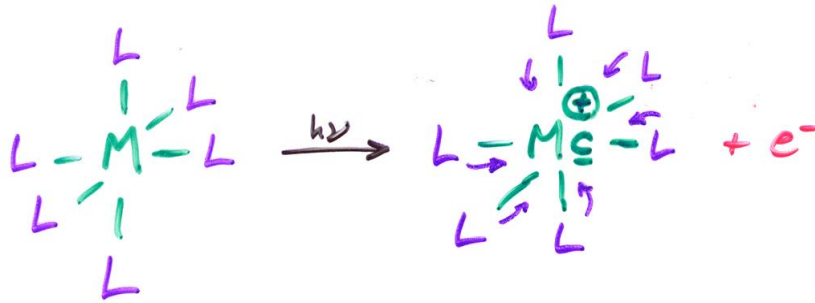
$T_\sigma$  in some papers and of opposite sign to  $T_\pi$  in some papers. Just a sign convention in the codes.

Interaction parameters describing the charge transfer now have to be input. These are different for each case, and are mostly empirically derived by fitting to XAS a/o XPS spectra.

$T_\pi$  in some papers and of opposite sign to  $T_\sigma$  in some papers

# Localized configuration interaction approach to spectrum simulation: Anderson impurity model for PS, XAS, XES

(SUGANO, LARSSON → SAWATZKY, VANDER LAAN,  
FUJIMORI, OH, ET AL.)



$\underline{c}$  = CORE HOLE ON METAL

$\bar{L}$  = VALENCE (?) HOLE ON LIGAND

$$\psi_i = a_0 |d^n\rangle + \sum_m a_m |d^{(n+m)} \bar{L}^m\rangle$$

$$\psi_f = b_0 |\underline{c} d^n\rangle + \sum_m b_m |\underline{c} d^{(n+m)} \bar{L}^m\rangle$$

WITH INTERACTIONS OF:

$10Dq$  = CRYSTAL FIELD (OFTEN NEGLECTED)

$\Delta$  = LIGAND-TO-METAL CHARGE TRANSF. ENERGY  
 $= E(d^{n+1} \bar{L}) - E(d^n)$

$U_{pd} = U$  = d-d COULOMB REPULSION ENERGY  
 $= E(d^{n-1}) + E(d^{n+1}) - 2E(d^n)$

$T$  = LIGAND p-TO-METAL d HYBRIDIZATION  
 $= \langle d_\alpha | \hat{H} | p_\alpha \rangle$  ( $\alpha$  = SAME SYMMETRY)

$U_{dd} = Q$  = CORE-HOLE-TO-d INTERACTION:  $\langle \underline{c} | \hat{H} | d \rangle \approx J_{cd}$  = coulomb integral

Good discussion of model:  
Bocquet & Fujimori, J. Elect.  
Spect. & Rel. Phen. 82, 87  
(1996)

$$\rho(e_k) = \sum_f |\langle \Psi_f | c | \Psi_g \rangle|^2 \delta(h\nu - e_k - E_f)$$

By now:

CTM4XAS program  
for calculating this

for some cases:

[http://www.anorg.chem.uu.nl/  
people/staff/FrankdeGroot/  
multiplet1.htm](http://www.anorg.chem.uu.nl/people/staff/FrankdeGroot/multiplet1.htm)

WITH INTENSITIES FROM SUDDEN APPROX.

AS:

$$I(E_{kin}) \propto \sum_{f,k} |\langle \Psi_f(N-1, k) | \Psi_g(N-1, k) \rangle|^2$$

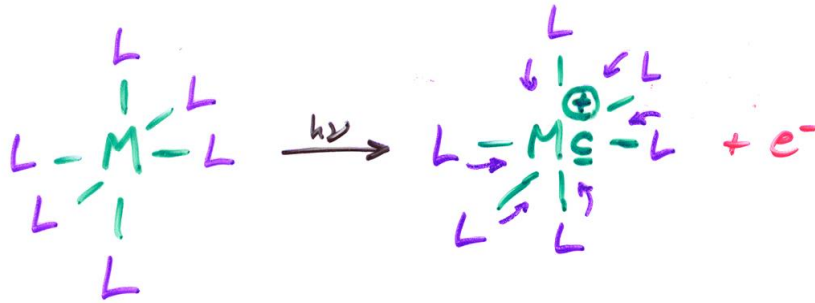
$$\delta(h\nu - E_f - E_{kin})$$

WHERE:  $\Psi_g(N-1, k) = \psi_i$  (N WITH k HOLE =  $\underline{c}$ )

$\underline{c}$  = CORE HOLE

# Localized configuration interaction approach to spectrum simulation: Anderson impurity model for PS, XAS, XES

(SUGANO, LARSSON → SAWATZKY, VANDER LAAN,  
FUJIMORI, OH, ET AL.)



$\underline{c}$  = CORE HOLE ON METAL  
 $\underline{v}$  = VALENCE (?) HOLE ON LIGAND

$$\psi_i = a_0 |d^n\rangle + \sum_m a_m |d^{(n+m)} \underline{v}^m\rangle$$

$$\psi_f = b_0 |\underline{c} d^n\rangle + \sum_m b_m |\underline{c} d^{(n+m)} \underline{v}^m\rangle$$

WITH INTERACTIONS OF:

$10Dq$  = CRYSTAL FIELD (OFTEN NEGLECTED)

$\Delta$  = LIGAND-TO-METAL CHARGE TRANSF. ENERGY  
=  $E(d^{n+1} \underline{v}) - E(d^n)$

$U_{pd} = U$  = d-d COULOMB REPULSION ENERGY  
=  $E(d^{n-1}) + E(d^{n+1}) - 2E(d^n)$

$T$  = LIGAND p-TO-METAL d HYBRIDIZATION  
=  $\langle d_\alpha | \hat{H} | p_\alpha \rangle$  ( $\alpha$  = SAME SYMMETRY)

$U_{dd} = Q$  = CORE-HOLE-TO-d INTERACTION:  $\langle \underline{c} | \hat{H} | d \rangle \approx J_{cd}$  = coulomb integral

Good discussion of model:  
Bocquet & Fujimori, J. Elect.  
Spect. & Rel. Phen. 82, 87  
(1996)

$$\rho(e_k) = \sum_f |\langle \Psi_f | c | \Psi_g \rangle|^2 \delta(h\nu - e_k - E_f)$$

By now:

CTM4XAS program  
for calculating this

for some cases:

[http://www.anorg.chem.uu.nl/  
people/staff/FrankdeGroot/  
multiplet1.htm](http://www.anorg.chem.uu.nl/people/staff/FrankdeGroot/multiplet1.htm)

WITH INTENSITIES FROM SUDDEN APPROX.

AS:

$$I(E_{kin}) \propto \sum_{f,k} |\langle \Psi_f(N-1, k) | \Psi_g(N-1, k) \rangle|^2$$

$$\delta(h\nu - E_f - E_{kin})$$

WHERE:  $\Psi_g(N-1, k) = \psi_i$  (N WITH k HOLE =  $\underline{c}$ )

$\underline{c} = \underline{c}$  = CORE HOLE

**From Bocquet & Fujimori, J. Elect.Spect. & Rel. Phen. 82, 87 (1996):**

The electronic structures of transition-metal (TM) compounds, particularly the 3d TM halides, oxides and chalcogenides, have long provided intriguing problems for physicists and chemists. These compounds are highly correlated electron systems where the essential physics can be described in terms of a few interaction strengths, namely the on-site d–d coulombic repulsion energy  $U$ , the ligand-to-metal charge-transfer energy  $\Delta$ , and the ligand p–metal d hybridization strength  $T$ . Core-level X-ray photoemission spectroscopy (XPS) is a useful probe of the valence electronic structures of TM compounds, and has been successfully used in recent years to extract parameter values for these interaction strengths [1–5]. In

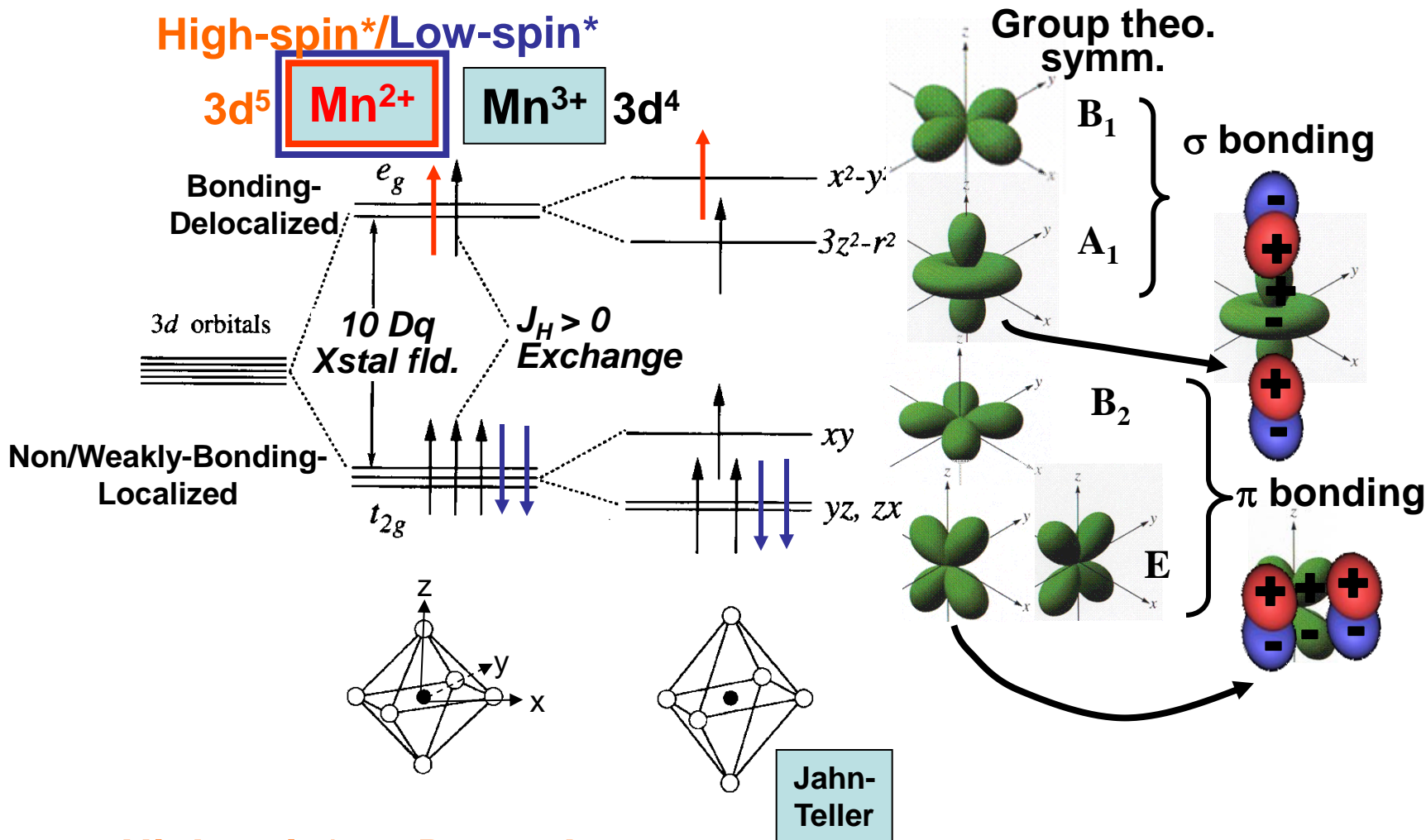
**From CTM manual:**

**C: Charge Transfer Parameters**

- **Delta:** This is the charge transfer parameter  $\Delta$ , which gives the energy difference between the (centers of the)  $3d^N$  and  $3d^{N+1}\underline{L}$  configurations. The effective value of  $\Delta$  ( $\Delta_{\text{eff}}$ ) is affected by the multiplet and crystal field effects on each configuration. In the next version, the value of  $\Delta_{\text{eff}}$  will be given in a parameter-output file.
- **Udd:** This is the value of the Hubbard  $U$ .
- **Upd:** This is the core hole potential. In case of XAS spectra, only the difference between  $U_{\text{pd}}$  and  $U_{\text{dd}}$  is important.
- **Hopping T:** The hopping parameters are given for the 4 symmetries in tetragonal symmetry  $A_1$ ,  $B_1$ ,  $E$  and  $B_2$ .  $A_1$  ( $z^2$ ) and  $B_1$  ( $x^2-y^2$ ) are part of the  $e_g$ -orbitals and  $E$  ( $xz$ ,  $yz$ ) and  $B_2$  ( $xy$ ) are part of the  $t_{2g}$  orbitals. In Oh symmetry the values of  $A_1=B_1$  and  $E=B_2$ . (This is not yet automatic in the test-version).

**Originated in the Hubbard Model: Ashcroft and Mermin, pp. 689-691**

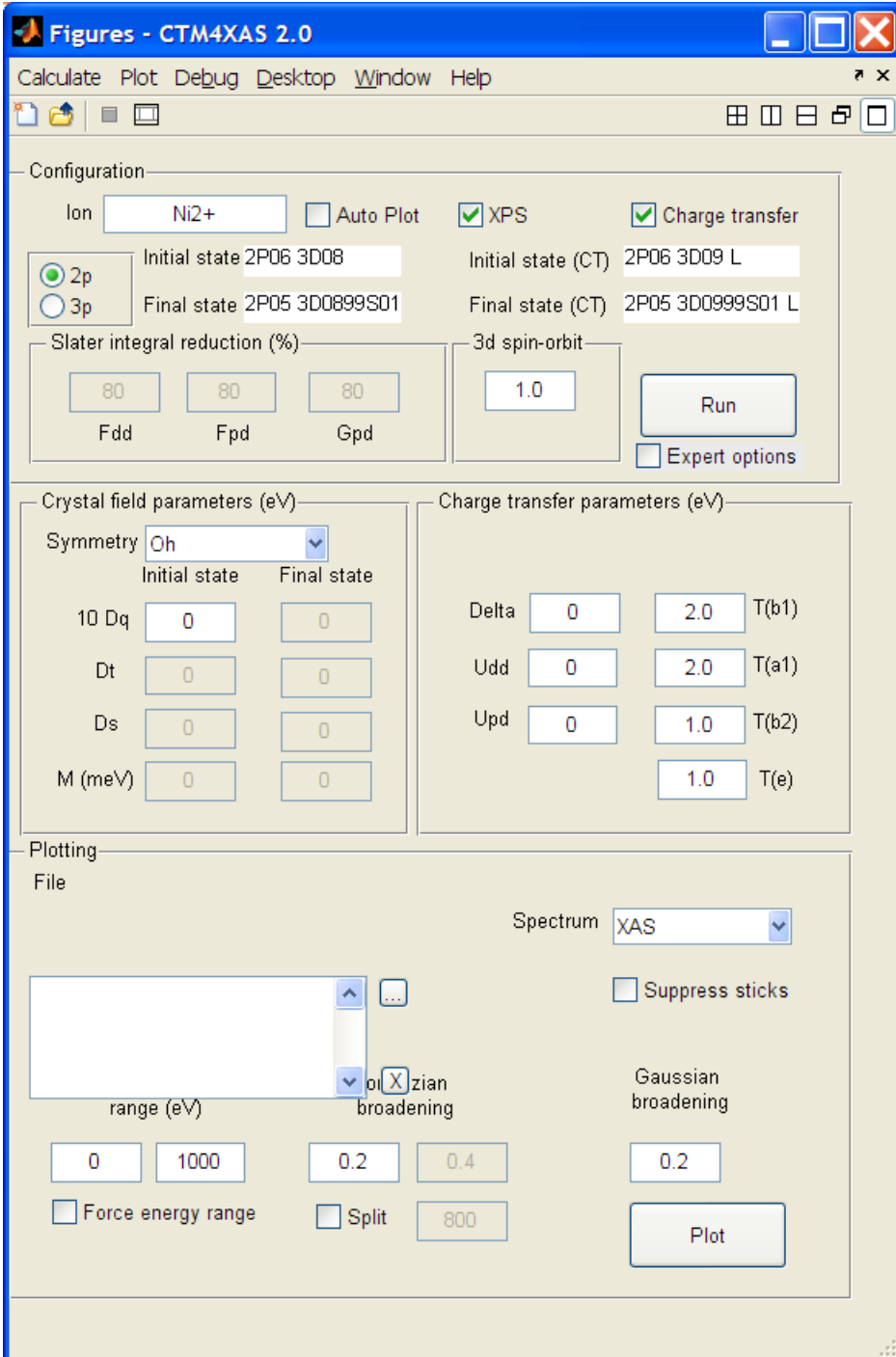
# E.g.—Crystal field in $Mn^{3+}$ & $Mn^{2+}$ with negative octahedral ligands



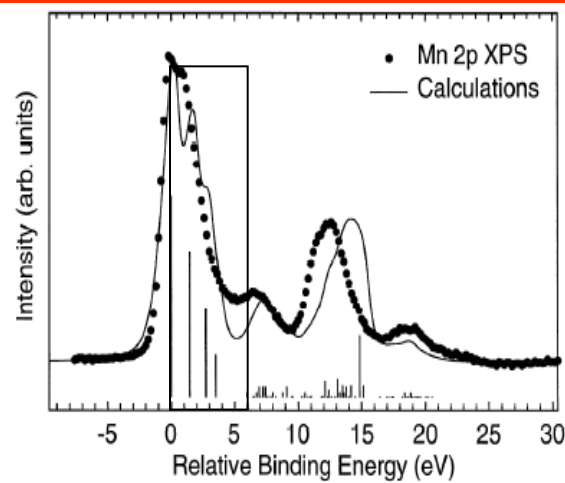
High-spin\*:  $10Dq \ll J_H$

Low-spin\*:  $10Dq \gg J_H$





Initial and final states in XPS have different numbers of electrons, and screened  $d^{n+1}$  configurations mixed in are now present. Notation is odd for the latter.



More recent data:  
 Phys. Rev. B 63, 115119 (2001)  
 In problem set

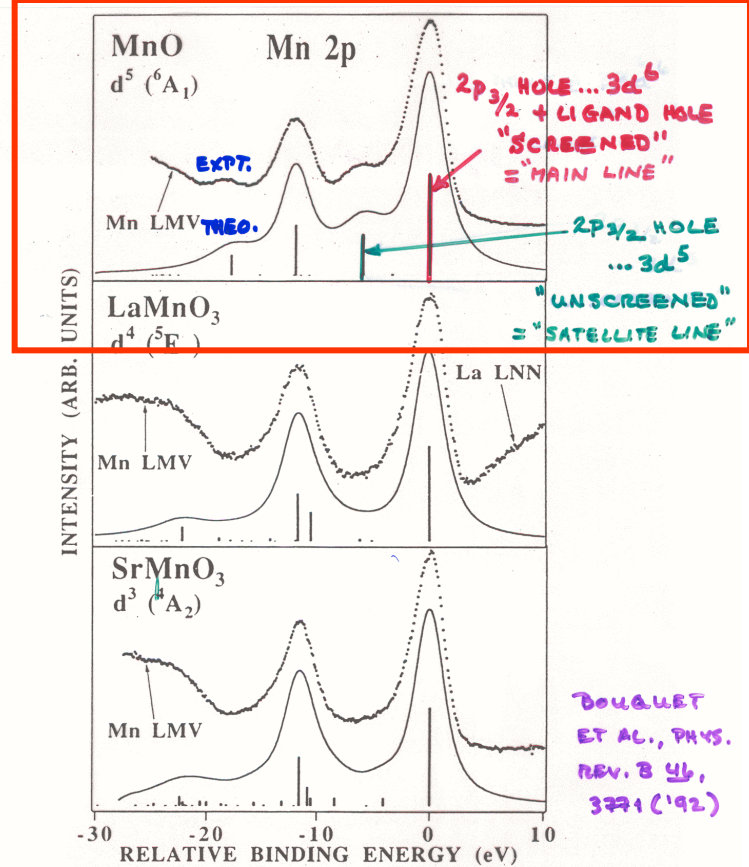


FIG. 1. Theoretical 2p core-level XPS spectra (solid line) compared with experimental data (dots) after background subtraction for Mn cations with varying valence. Emission due to the Mn LMV Auger peak is observed on the high-binding-energy side of the 2p<sub>1/2</sub> spin-orbit peak, partially obscuring the 2p<sub>1/2</sub> satellite structure.

For octahedral coord.:

$$T_{\sigma} = \sqrt{3}(pd\sigma), \quad T_{\pi} = 2(pd\pi)$$

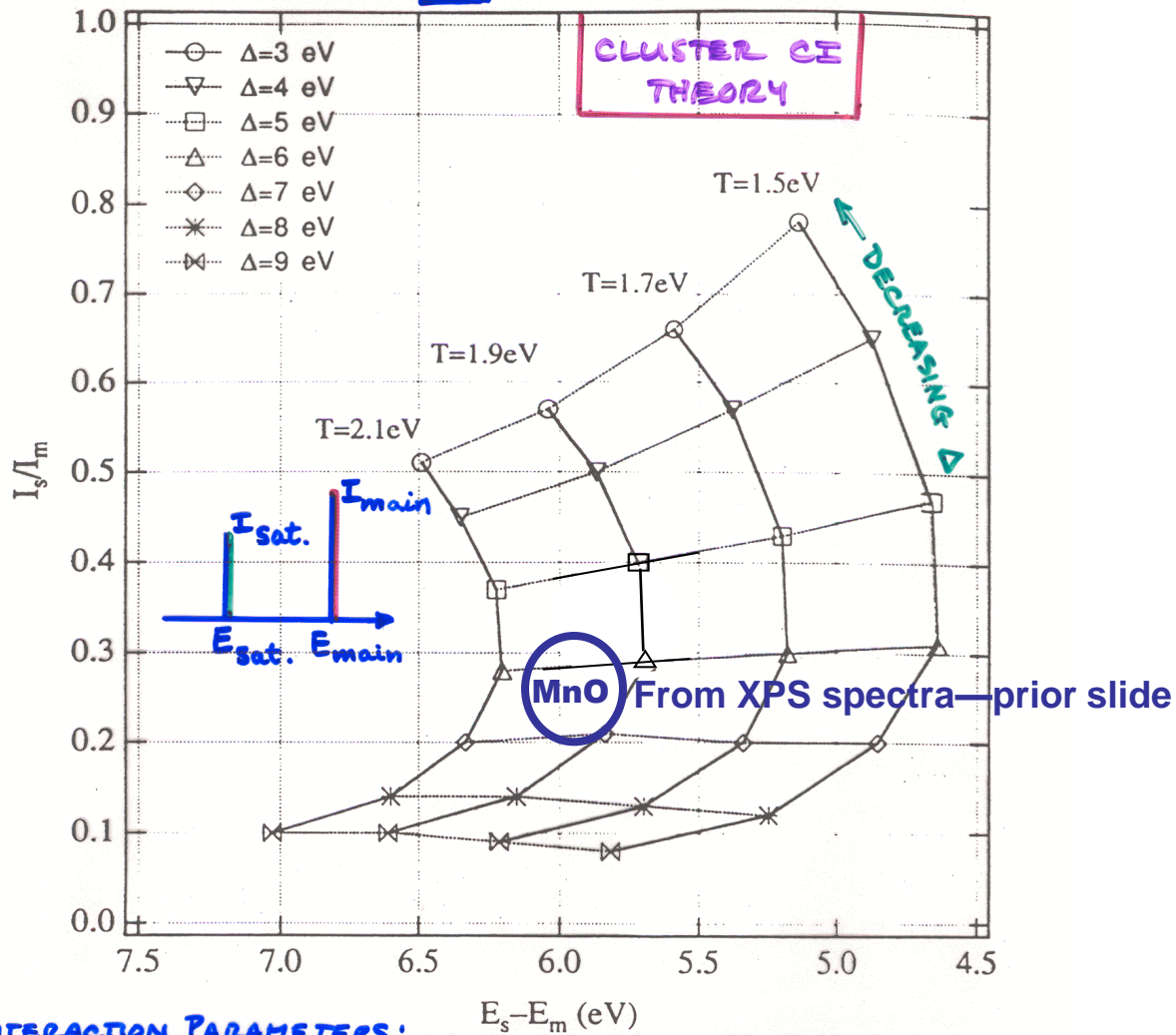
$$T_{\sigma} / T_{\pi} \propto (pd\sigma) / (pd\pi) \approx -2.2.$$

| Compound           | d <sup>n</sup> | Valence | Δ   | U   | (pdσ) | Δ <sub>eff</sub> | U <sub>eff</sub> | Main peak         | Satellite peak                 | Ref.      |
|--------------------|----------------|---------|-----|-----|-------|------------------|------------------|-------------------|--------------------------------|-----------|
| SrMnO <sub>3</sub> | d <sup>3</sup> | 4+      | 2.0 | 7.8 | -1.5  | -0.2             | 7.1              | d <sup>4</sup> L̄ | d <sup>4</sup> L̄              | This work |
| LaMnO <sub>3</sub> | d <sup>4</sup> | 3+      | 4.5 | 7.5 | -1.8  | 1.8              | 6.8              | d <sup>5</sup> L̄ | d <sup>6</sup> L̄ <sup>1</sup> | This work |
| MnO                | d <sup>5</sup> | 2+      | 6.5 | 7.0 | -1.1  | 8.8              | 11.6             | d <sup>6</sup> L̄ | d <sup>5</sup>                 | This work |
|                    |                |         | 7.0 | 7.5 | -0.9  |                  |                  |                   |                                | 3         |

From later paper (see next slides) 6.5    6.0    -1.99/√3 = -1.1    U/Q = 0.7-1.0, a best fit no. is 0.83.    10Dq is 1.0-2.0 eV

# ANALYSIS VIA ANDERSON IMPURITY MODEL

Mn<sup>2+</sup> (HS) U=6.0 eV



**INTERACTION PARAMETERS:**

$U_{dd} \equiv U = 3d-3d$  COULOMB REPULSION ENERGY

$\Delta =$  LIGAND-TO-METAL CHARGE TRANSFER ENERGY

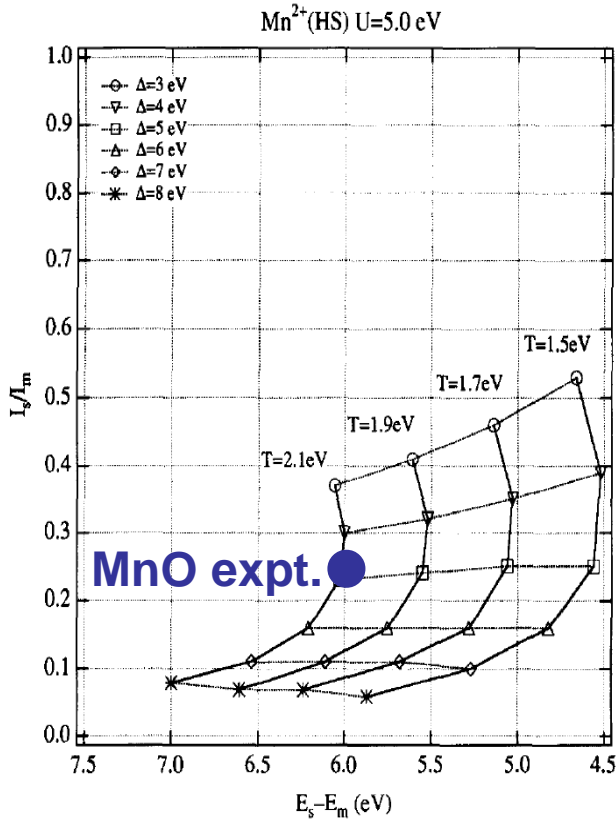
$T =$  LIGAND p - METAL 3d HYBRIDIZATION ENERGY

$U_{pd} \equiv Q =$  CORE HOLE-3d COULOMB

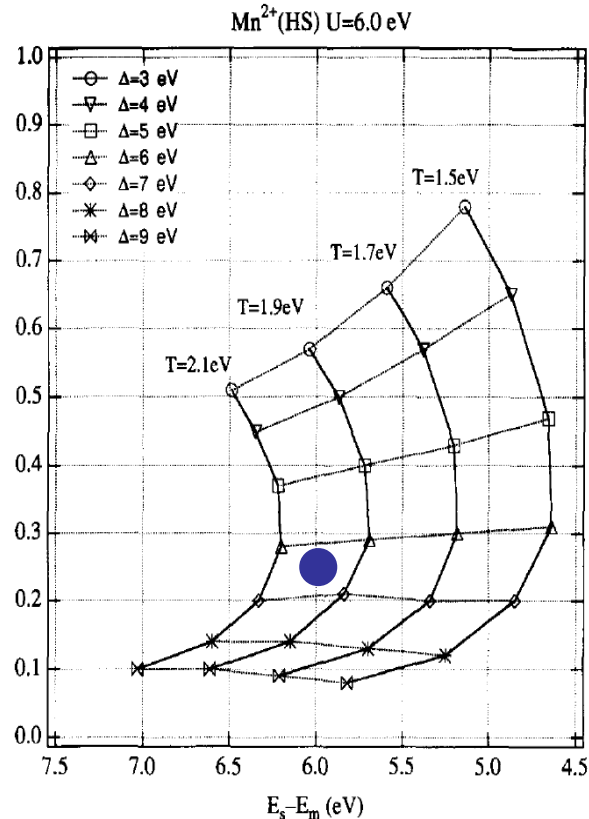
BOLQUET ET AL.,  
J. EL. SP. 82, 87 (196)

# Estimating the parameters from fits to the experimental satellite intensity and spacing

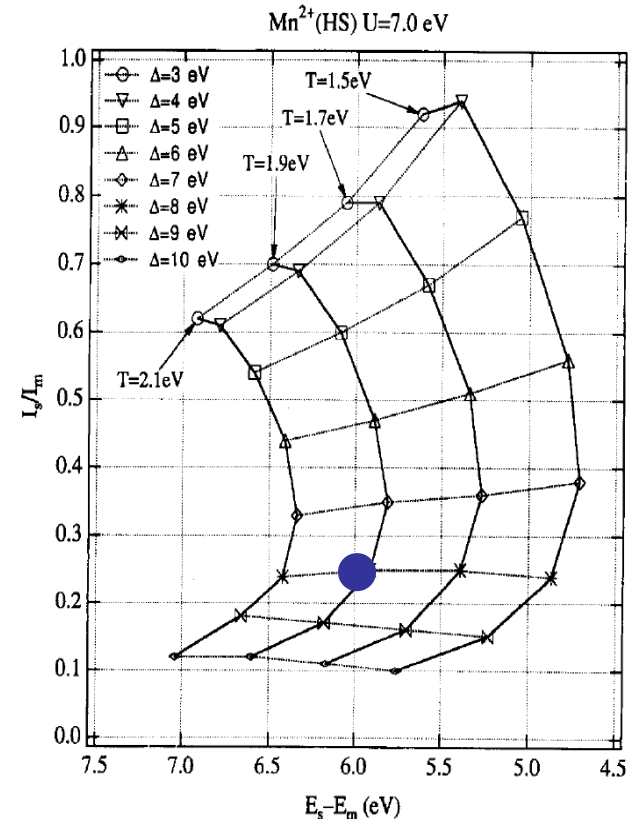
From Bocquet & Fujimori, J. Elect.Spect. & Rel. Phen. 82, 87 (1996)



**U = 5.0, T = 2.08, Δ = 4.8**



**U = 6.0, T = 1.99, Δ = 6.5**

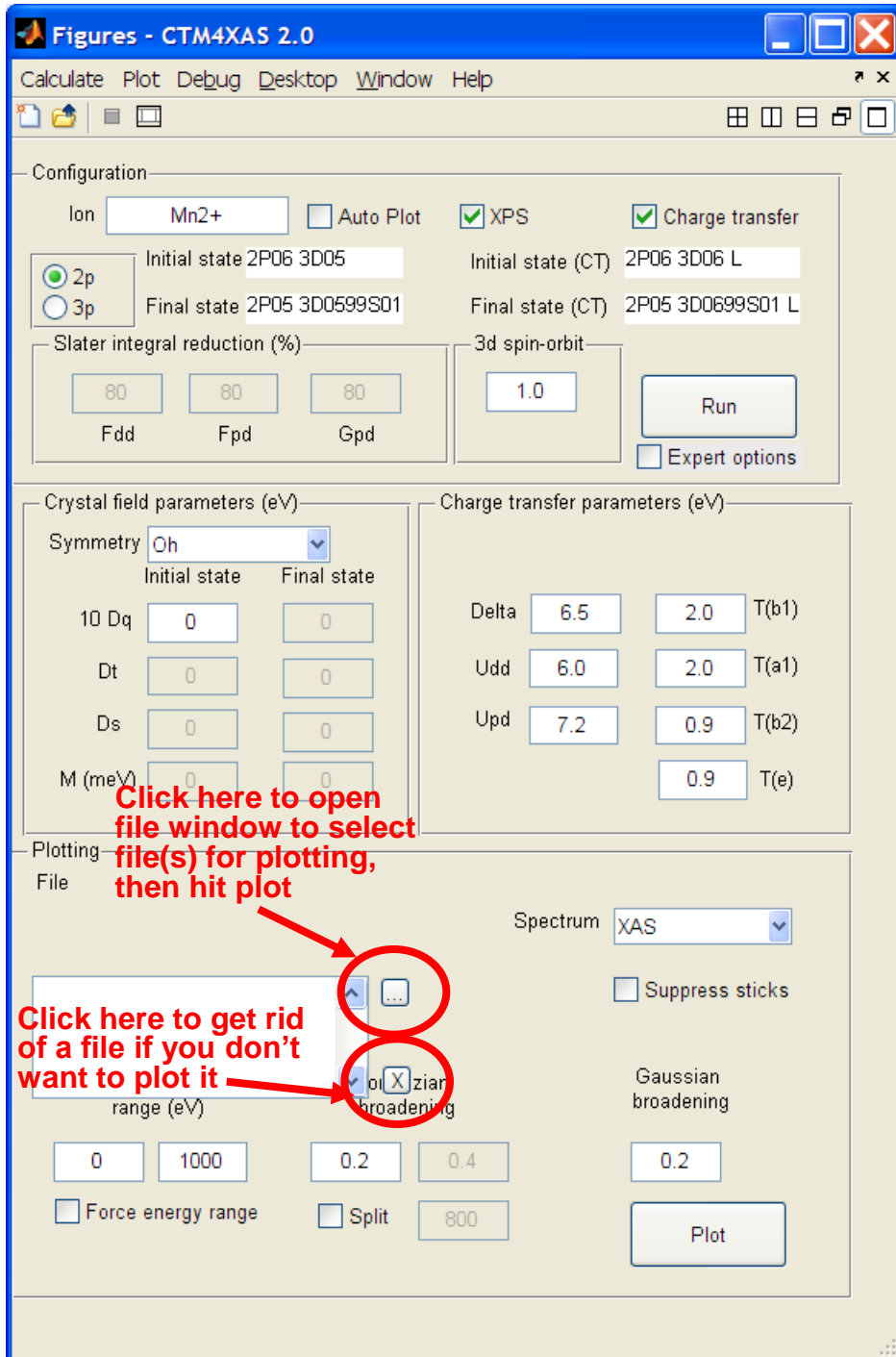


**U = 7.0, T = 1.93, Δ = 8.0**

Crystal field 10Dq assumed zero in this study,  
 but a value of about ~1.0-2.0 eV has been estimated from previous  
 optical, XAS, and EELS measurements:

Garvie et al., Physical Chemistry of Materials 21, 191 (1994)

Muller et al., Phys. Rev. B 78, 085438 (2008)



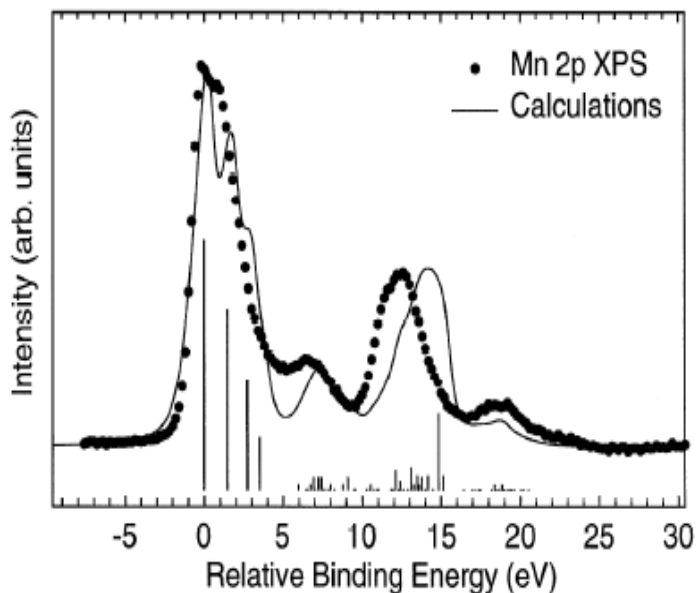
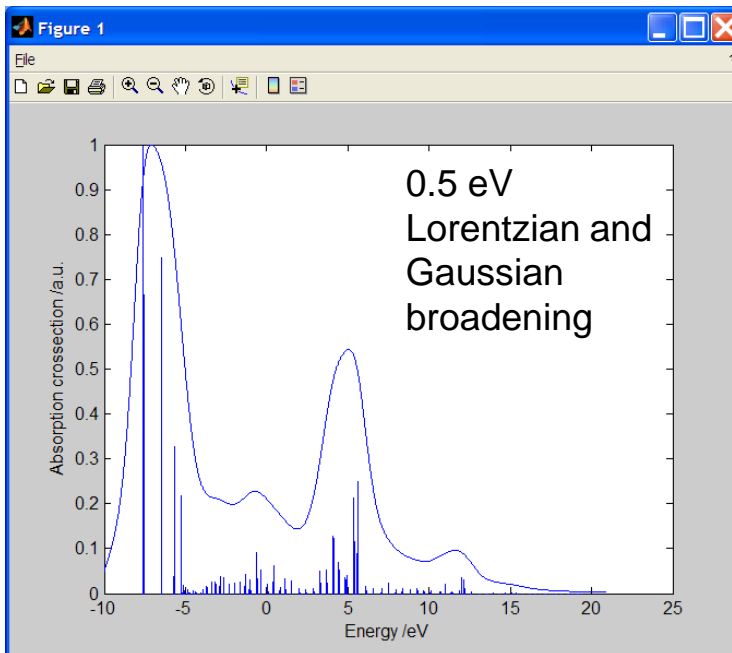
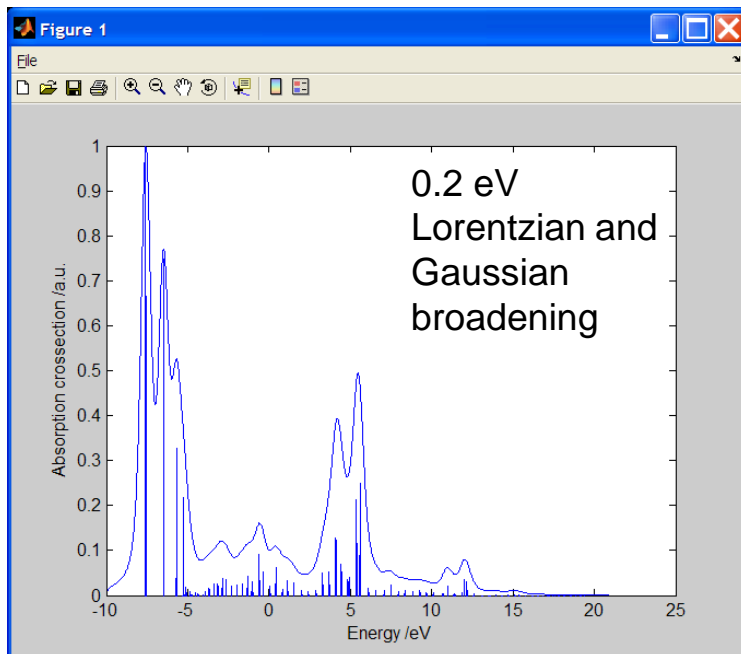
Here is a reasonable choice of parameters for  $Mn^{2+}$ , with no crystal field splitting

Hit run, wait until program asks for a file name for the output, input this and run will begin. Wait while program steps are executing in command window, example below. XPS is much slower than XAS.

```

C:\cowan\ctm4xas20\ctm4xas20.exe
looking for GROUND S1- 0+ S1- 758 758 Factor = 0.0000000
routine=ETR file= 15 line= 20785 col= 66 =>
REDUCEDMATRIX GROUND S1- 0+ S1- DT 758 758<<
looking for GROUND S1- 0+ S1- 758 758 Factor = 0.0000000
routine=ETR file= 15 line= 22055 col= 66 =>
REDUCEDMATRIX GROUND S1- 0+ S1- DS 758 758<<
looking for GROUND S1- 0+ S1- 758 758 Factor = 0.0000000
routine=ETR file= 15 line= 22975 col= 66 =>
REDUCEDMATRIX GROUND S1- 0+ S1- SPIN 758 758<<
looking for EXCITE S1- 0+ S1- 6302 6302 Factor = 1.0000000
routine=ETR file= 15 line= 23517 col= 68 =>
REDUCEDMATRIX EXCITE S1- 0+ S1- HAMILTONIAN 6302 6302<<
looking for EXCITE S1- 0+ S1- 6302 6302 Factor = 0.0000000
routine=ETR file= 15 line= 24581 col= 68 =>
REDUCEDMATRIX EXCITE S1- 0+ S1- 10D0 6302 6302<<
looking for EXCITE S1- 0+ S1- 6302 6302 Factor = 0.0000000
routine=ETR file= 15 line= 26943 col= 68 =>
REDUCEDMATRIX EXCITE S1- 0+ S1- DT 6302 6302<<
looking for EXCITE S1- 0+ S1- 6302 6302 Factor = 0.0000000
routine=ETR file= 15 line= 29369 col= 68 =>
REDUCEDMATRIX EXCITE S1- 0+ S1- DS 6302 6302<<
looking for EXCITE S1- 0+ S1- 6302 6302 Factor = 0.0000000
routine=ETR file= 15 line= 31031 col= 68 =>
REDUCEDMATRIX EXCITE S1- 0+ S1- SPIN 6302 6302<<
  
```

And finally you will get something like this, with no crystal field splitting:



Which compares pretty well with expt. and another similar calculation from the literature, esp. if theory is broadened more.

Could try adding crystal field, as next slide suggests.

Now just switching to XAS gives you:

